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# ON SOME RELATIONS CONNECTING FLUID DYNAMICS AND BI-COMPLEX ANALYSIS

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Abstract: In the paper our main target is to derive some results focusing some connection between fluid dynamics and bi-complex analysis which in fact is the most recent mathematical tool to develop the theory of complex analysis.

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## **1** INTRODUCTION, DEFINITIONS AND NOTATIONS.

The most recent advancement of the theory of complex numbers lies in the progress of bi-complex analysis. According to Segree [9], a bi-complex number  $\xi$  is defined as follows:

$$\xi = x_0 + i_1 x_1 + i_2 x_2 + i_1 i_2 x_3$$

where  $x_0, x_1, x_2$  and  $x_3$  are real numbers with  $i_1^2 = i_2^2 = -1$  and  $i_1 i_2 = i_2 i_1$ .

The set of all bi-complex numbers is generally denoted by  $C_2$ . In the theory of bi-complex numbers, the sets of real numbers and complex numbers are generally denoted by  $C_0$  and  $C_1$  respectively. Thus

 $C_2 = \{\xi : \xi = x_0 + i_1 x_1 + i_2 x_2 + i_1 i_2 x_3 ; x_0, x_1, x_2, x_3 \in C_0\}.$ 

Equivalently, we may write  $C_2 = \{\xi : \xi = z_1 + i_2 z_2 ; z_1 z_2 \in C_1\}.$ 

We have seen some illuminating works on the recent advancement of different aspect of bi-complex analysis in Michiji Futagawa [2], E. Hille [3], D. Riley [4], G. Baley Price [1]. In the present paper we would like to establish some results in fluid dynamics in close- connection to bi-complex analysis. In fact, the paper is an improved version of Datta and Sen [15] and therefore all the preliminary theories and definitions on bi-complex analysis as required here are omitted.

Now, let us define a function as follows:

Let  $\Psi: [0, \infty) \to (0, \infty)$  be a non-decreasing unbounded function, satisfying the following two conditions:

(*i*) 
$$\lim_{r \to \infty} \frac{\log^{[p]}(r)}{\log^{[q]}[\Psi(r)]} = 1$$

and

$$(ii)\lim_{r\to\infty}\frac{\log^{[q]}(\alpha r)}{\log^{[q]}\Psi[(r)]}=1,$$

for some  $\alpha > 1$  and *p*, *q* are any two posotive integers.

With the help of function  $\Psi$  as defined earlier the following definitions may be given :

**Definition 1.** The  $\Psi$  –order  $\rho_{(F,\Psi)}$  of a bicomplex meromorphic function  $F(w) = F_{e_1}(z_1 - i_1 z_2)e_1 + (z_1 + i_1 z_2)e_2$ 

$$\rho_{(F,\Psi)} = max\left\{\rho_{(F_{e_1,\Psi})}, \rho_{(F_{e_2,\Psi})}\right\}$$

where

is defined as

$$\rho_{F_{e_i},\Psi} = \limsup_{r_i \to \infty} \frac{\log \log M_i(r_i, F_{e_i})}{\log[\Psi(r_i)]} for \ i = 1,2$$

**Remark 1.** The  $\Psi$  –lower order  $\lambda_{(F,\Psi)}$  of a bicomplex meromorphic function is defined as

$$\lambda_{(F,\Psi)} = \min\left\{\lambda_{(F_{e_1,\Psi})}, \lambda_{(F_{e_2,\Psi})}\right\},\,$$

where

$$\lambda_{(F_{e_i},\Psi)} = \liminf_{r_i \to \infty} \frac{\log \log M_i(r_i, F_{e_i})}{\log [\Psi(r_i)]} \text{ for } i = 1,2.$$

**Remark 2.** The  $\Psi$  -hyper order  $\bar{\rho}_{(F,\Psi)}(\Psi$  -hyper lower order  $\bar{\lambda}_{(F,\Psi)})$  and the generalized  $\Psi$  -order  $\rho_{(F,\Psi)}^{(k)}$  (generalized  $\Psi$  -lower order  $\lambda_{(F,\Psi)}^{(k)}$ ) can also be defined in a similar way.

**Definition 2.** The  $\Psi$  -type of F  $\sigma_{(F,\Psi)}$  of a bicomplex meromorphic function is defined as

$$\sigma_{(F,\Psi)} = max\left\{\sigma_{(F_{e_1,\Psi})}, \sigma_{(F_{e_2,\Psi})}\right\}$$

where

$$\sigma_{(F_{e_i},\Psi)} = \text{limsup}_{r_i \to \infty} \frac{\log M_i(r_i, F_{e_i}, \Psi)}{\Psi(r_i^{\sigma(F_{e_i},\Psi)})} \text{ and } 0 < \rho_{(F_{e_i},\Psi)} < \infty \text{ for } i = 1,2$$

**Definition 3.** Let F(w) be an entire function of  $\Psi$  –order zero. Then the quantities  $\rho_{(F,\Psi)}^*$  and  $\lambda_{(F,\Psi)}^*$  can be defined as

$$\rho_{(F,\Psi)}^{*} = max \left\{ \rho_{(F_{e_{1}},\Psi)}^{*}, \rho_{(F_{e_{2}},\Psi)}^{*} \right\}$$
and
$$\lambda_{(F,\Psi)}^{*} = min \left\{ \lambda_{(F_{e_{1}},\Psi)}^{*}, \lambda_{(F_{e_{2}},\Psi)}^{*} \right\}$$
where
$$\rho_{(F_{e_{i},\Psi)}}^{*} = \limsup_{r_{i} \to \infty} \frac{loglogM_{i}(r_{i}, F_{e_{i}})}{loglog[\Psi(r_{i})]}$$
and
$$\lambda_{F_{e_{i}}}^{*} = \liminf_{r_{i} \to \infty} \frac{loglogM_{i}(r_{i}, F_{e_{i}})}{loglog[\Psi(r_{i})]} for i = 1,2.$$

**Definition 4.**Let F(w) be an entire function order zero. Then the quantities  $\rho_{(F,\Psi)}^{**}$  and  $\lambda_{(F,\Psi)}^{**}$  can be defined as

$$\rho_{(F,\Psi)}^{**} = max \left\{ \rho_{(F_{e_1},\Psi)}^{**}, \rho_{(F_{e_2},\Psi)}^{**} \right\} and$$

$$\lambda_{(F,\Psi)}^{**} = min \left\{ \lambda_{(F_{e_1},\Psi)}^{**}, \lambda_{(F_{e_2},\Psi)}^{**} \right\} where$$

$$\rho_{F_{e_i}}^{**} = \limsup_{r_i \to \infty} \frac{\log M_i(r_i, F_{e_i})}{\log [\Psi(r_i)]} and \lambda_{F_{e_i}}^{**} = \liminf_{r_i \to \infty} \frac{\log M_i(r_i, F_{e_i})}{\log [\Psi(r_i)]} for \ i = 1, 2.$$

**Definition 5.** (Factorization of F(w)) Let F(w) be a bicomplex meromorphic function on  $T \subset C_2$ . Then F is said to have f and g as left and right factors respectively if Fei has fei and gei as left and right factors respectively for i=1,2, i.e., f<sub>ei</sub> is meromorphic and g<sub>ei</sub> is entire for i=1,2.

**Definition 6.** (Complex potential flow) If  $f(z) = u(x, y) + iv(x, y) \in C_1$  be a complex function where  $u(x, y) \in R^2$  and  $v(x, y) \in R^2$  satisfy the Cauchy-Riemann equations, i.e.,  $\frac{\partial u}{\partial x} = \frac{\partial v}{\partial y}$ ,  $\frac{\partial u}{\partial y} = -\frac{\partial v}{\partial x}$  and Laplace's equation, i.e.,  $\frac{\partial_u^2}{\partial x^2} + \frac{\partial_u^2}{\partial y^2} = 0$ , then f(z) can be termed as a complex potential fluid flow.

**Definition 7.**(Bicomplex potential fluid flow) If  $f(w) = f(z_1 + i_2 z_2) = f_{e_1}(z_1 - i_1 z_2)e_1 + f_{e_2}(z_1 + i_1 z_2 e_2)e_1$  be the idempotent composition of two complex functions, with  $f_{e_1}(z_1 - i_1 z_2) = u(z_1, z_2) - i_1 v(z_1 - z_2) \in C_1$  and  $f_{e_2}(z_1 - i_1 z_2) = u(z_1, z_2) + i_1 v(z_1, z_2) \in C_1$  where  $u(z_1, z_2)$  and  $v(z_1, z_2)$  satisfy Cauchy-Riemann equations and Laplace's equation, i.e.,  $\frac{\partial u}{\partial z_1} = \frac{\partial v}{\partial z_2}, \frac{\partial u}{\partial z_2} = \frac{\partial v}{\partial z_2}$  $-\frac{\partial v}{\partial z_2} \text{ and } \frac{\partial_u^2}{\partial z_1^2} + \frac{\partial_v^2}{\partial z_2^2} = 0, \\ \frac{\partial_u^2}{\partial z_1^2} + \frac{\partial_v^2}{\partial z_2^2} = 0, \text{ therefore } f_{e_1}(z_1 - i_1 z_2) \text{ and } f_{e_2}(z_1 + i_1 z_2) \text{ can be termed as complex}$ potential fluid flows. So, f(w) can be termed as a composition of two different potential fluid flows  $f_{e_1}$  and  $\begin{array}{c} f_{e_2} \\ \mathbf{3} \end{array} \mathbf{LEMMA.} \end{array}$ 

In this section we present some lemmas which will be needed in the sequel.

**Lemma 1 [8] [15].** If f(z) = u(x, y) + iv(x, y) be complex potential fluid flow defined in the region  $\{y > 0\}$  satisfying the following properties :

- (i) f(z) is continuously differentiable in the region  $\{y \ge 0\}$ ,
- f'(z) is parallel to the x-axis when y = 0 and (ii)
- f'(z) is uniformly bounded in  $\{y > 0\}$ , then the order and lower order of f(z) are zero. (iii)

**Corollary 1.** If  $f(w) = f(z_1 + i_2 z_2) = f_{e_1}(z_1 - i_1 z_2)e_1 + f_{e_2}(z_1 + i_1 z_2)e_2$  be an idempotent composition of two complex potential fluid flows satisfying the following properties :

- (i)
- (ii)
- $f_{e_1}$  and  $f_{e_2}$  are continuously differentiable in the region  $\{y \ge 0\}$  $f_{e_1}'$  and  $f_{e_2}'$  are parallel to the *x*-axis when y = 0 and  $f_{e_1}'$  and  $f_{e_2}'$  are uniformly bounded in  $\{y > 0\}$ , then the order and lower order of f(w) are zero. (iii)

**Lemma 2** [14]. If f(z) and g(z) are any two entire functions, then for all sufficiently large values of r,

$$M\left(\frac{1}{8}M\left(\frac{r}{2},g\right) - |g(0),f|\right) \le M(r,fog) \le M(M(r,g),f).$$

Lemma 3[14]. If f be entire and g be a transcendental entire function of finite lower order, then for any  $\delta > 0$ ,

$$M(r^{1+\delta}, fog) \ge M(M(r, g), f). \qquad (r \ge r_0)$$

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Lemma 4 [13]. If F has f and gas left and right factors, then we always have the following factorization :

$$F(w) = f(g(w)).$$

### THEOREMS. 4

In this section we present our main results of the paper.

**Theorem 1.** If  $f(w) = f(z_1 + i_2 z_2) = f_{e_1}(z_1 - i_1 z_2)e_1 + f_{e_2}(z_1 + i_1 z_2)e_2$  be an idempotent composition of two complex potential fluid flows  $f_{e_1}$  and  $f_{e_2}$  satisfying the following properties :

- $f_{e_1}$  and  $f_{e_2}$  are continuously differentiable in the region  $\{y \ge 0\}$ (i)
- $f_{e_1}^{\prime}$  and  $f_{e_2}^{\prime}$  are parallel to x-axis when y = 0 and  $f_{e_1}^{\prime}$  and  $f_{e_2}^{\prime}$  are uniformly bounded in  $\{y > 0\}$ , (ii)
- (iii)

then 
$$ho_{f}^{'}=1$$
 and  $\lambda_{f}^{'}=1.$ 

**Proof.** From the definitions of  $\rho_{f,\psi}^{**}$  and  $\lambda_{f,\psi}^{**}$  and using Definition 4, we have for arbitrary positive  $\varepsilon_1$ ,  $\varepsilon_2$  and all sufficiently large values of  $r_1, r_2$ 

$$\begin{split} \log M_1(r_1, f_{e_1}) &\leq \left(\rho_{(f_{e_1}, \Psi)}^{**} + \varepsilon_1\right) \log [\Psi(r_1)] \text{and} \\ \log M_2(r_2, f_{e_2}) &\leq \left(\rho_{(f_{e_2}, \Psi)}^{**} + \varepsilon_2\right) \log [\Psi(r_2)] \\ \text{Therefore, } \log \log M_1(r_1, f_{e_1}) &\leq \log \log [\Psi(r_1)] + O(1) \\ \text{and} \quad \log \log M_2(r_2, f_{e_2}) &\leq \log \log [\Psi(r_2)] + O(1) \\ \text{i.e.,} \frac{\log \log M_1(r_1, f_{e_1})}{\log \log [\Psi(r_1)]} &\leq \frac{\log \log [\Psi(r_1)] + O(1)}{\log \log [\Psi(r_1)]} \text{ and} \\ \frac{\log \log M_2(r_2, f_{e_2})}{\log \log r_2} &\leq \frac{\log \log [\Psi(r_2)] + O(1)}{\log \log [\Psi(r_2)]} \\ \text{i.e.,} \lim_{r_1 \to \infty} \frac{\log \log M_1(r_1, f_{e_1})}{\log \log [\Psi(r_1)]} &\leq 1 \text{ and} \lim_{r_2 \to \infty} \frac{\log \log M_2(r_2, f_{e_2})}{\log \log [\Psi(r_2)]} \leq 1 \end{split}$$

i.e.,  $\rho_{f_{e_1},\psi}^* \leq 1$  and  $\rho_{f_{e_2},\psi}^* \leq 1$  and therefore using Definition 3,

we have 
$$\rho_{f,\Psi}^* \le 1.$$
 (1)

Similarly, proceeding as above and using Definition 3,

we have 
$$\lambda_{f,\Psi}^* \leq 1.$$
 (2)

Again, for arbitrary positive  $\varepsilon_1, \varepsilon_2$  and all sufficiently large values of  $r_1, r_2$  we have

$$log M_1(r_1, f_{e_1}) \ge \left(\lambda_{(f_{e_1}, \Psi)}^{**} - \varepsilon_1\right) \log[\Psi(r_1)]$$
  
and  $log M_2(r_2, f_{e_2}) \ge \left(\lambda_{(f_{e_2}, \Psi)}^{**} - \varepsilon_2\right) log[\Psi(r_2)]$   
Therefore,  $log log M_1(r_1, f_{e_1}) \ge log log[\Psi(r_1)] + O(1)$ 

and  $log log M_2(r_2, f_{e_2}) \ge log log [\Psi(r_2)] + O(1)$ 

i.e., 
$$\frac{\log \log M_1(r_1, f_{e_1})}{\log \log [\Psi(r_1)]} \ge \frac{\log \log [\Psi(r_1)] + O(1)}{\log \log [\Psi(r_1)]}$$

and

$$\frac{\log\log M_2(r_2, f_{e_2})}{\log\log[\Psi(r_2)]} \ge \frac{\log\log[\Psi(r_2)] + O(1)}{\log\log[\Psi(r_2)]}$$
$$i.e., \limsup_{r_1 \to \infty} \frac{\log\log M_1(r_1, f_{e_1})}{\log\log[\Psi(r_1)]} \ge 1$$
$$and \limsup_{r_2 \to \infty} \frac{\log\log M_2(r_2, f_{e_2})}{\log\log[\Psi(r_2)]} \ge 1$$

i.e.,  $\rho^*_{(f_{e_1}, \Psi)} \ge 1$  and  $\rho^*_{(f_{e_2}, \Psi)} \ge 1$  and therefore using Definition 3,

we have 
$$\rho_{f,\Psi}^* \ge 1.$$
 (3)

Similarly, proceeding as above and using Definition 3,

we have 
$$\lambda_{f,\Psi}^* \ge 1.$$
 (4)

From (1) and (3) we have  $\rho_{f,\Psi}^* = 1$  and from (2) and (4) we have  $\lambda_{f,\Psi}^* = 1$ .

Thus the theorem follows.

**Example 1.** Let  $f(w) = w = z_1 + i_2 z_2 = (z_1 - i_1 z_2)e_1 + (z_1 + i_1 z_2)e_2 \in C_2$  (bicomplex space) be a bicomplex potential fluid flow in C<sub>2</sub>.

Therefore,  $f_{e_1} = z_1 - i_1 z_2 \in C_1$  (or *C*) and  $f_{e_2} = z_1 + i_1 z_2 \in C_1$ (or *C*)

Therefore, 
$$\rho_{f_{e_1},\Psi} = \limsup_{r_1 \to \infty} \frac{\log \log M_1(r_1, f_{e_1})}{\log [\Psi(r_1)]} = 0 \text{ as } M_1(r_1, f_{e_1}) \le |z_1 - i_1 z_2| \le r_1 + r_2$$
  
and  $\rho_{(f_{e_2},\Psi)} = \limsup_{r_2 \to \infty} \frac{\log \log M_2(r_2, f_{e_2})}{\log [\Psi(r_2)]} = 0 \text{ as } M_2(r_2, f_{e_2}) \le |z_1 + i_1 z_2| \le r_1 + r_2$ 

Hence  $\rho_{f,\Psi}^* = 0$  and similarly  $\lambda_{f,\Psi}^* = 0$ .

$$\rho_{(f_{e_1},\Psi)}^* = \limsup_{r_1 \to \infty} \frac{\log \log M_1(r_1, f_{e_1})}{\log \log [\Psi(r_1)]} = \limsup_{r_2 \to \infty} \frac{\log \log (r_1 + r_2)}{\log \log [\Psi(r_1)]} = 1, r_1 \text{ is fixed.}$$
  
And  $\rho_{(f_{e_2},\Psi)}^* = \limsup_{r_2 \to \infty} \frac{\log \log M_2(r_2, f_{e_2})}{\log \log [\Psi(r_2)]} = \limsup_{r_2 \to \infty} \frac{\log \log (r_1 + r_2)}{\log \log [\Psi(r_2)]} = 1, r_2 \text{ is fixed.}$ 

Therefore,  $\rho_{f,\Psi}^* = 1$  and similarly  $\lambda_{f,\Psi}^* = 1$ .

Similarly, we can also show that  $\rho_{f,\Psi}^{**} = 1$  and  $\lambda_{f,\Psi}^{**} = 1$ .

**Theorem 2.** Let f(w) and g(w) be any two bicomplex potential fluid flows satisfying the following properties :

- $f_{e_1}$ ,  $f_{e_2}$  are continuously differentiable in the region  $\{y \ge 0\}$ (i)
- (ii)
- $f_{e_1}^{'}, f_{e_2}^{'}$  are parallel to the *x*-axis when y = 0 $f_{e_1}^{'}, f_{e_2}^{'}$  are uniformly bounded in  $\{y > 0\}$ , such that  $\rho_f = 0$  and  $\lambda_f < \infty$ . (iii)

Then  $\rho_{f,\Psi} = \rho_{g,\Psi}$ .

**Proof.** Using Lemma 4, we can say that F(w) can be factorized to f(g(w)). Now, using Lemma 2 and Theorem 1 we have

$$\begin{split} \rho_{(F_{e_1},\Psi)} &= \rho_{\{(f_{e_1} \circ g_{e_1}),\Psi\}} = \limsup_{r_1 \to \infty} \frac{\log \log M_1(r_1, f_{e_1} \circ g_{e_1})}{\log [\Psi(r_1)]} \\ &\leq \limsup_{r_1 \to \infty} \frac{\log \log M(M(r_1, g_{e_1}), f_{e_1})}{\log \log M(r_1, g)} \cdot \limsup_{r_1 \to \infty} \frac{\log \log M(r_1, g_{e_1})}{\log [\Psi(r_1)]} \\ &\leq \rho_{(f_{e_1},\Psi)}^* \cdot \rho_{(g_{e_1},\Psi)} = 1 \cdot \rho_{(g_{e_1},\Psi)} = \rho_{g_{(e_1,\Psi)}} \end{split}$$

Similarly,

$$\rho_{(F_{e_2},\Psi)} \leq \rho_{(g_{e_2},\Psi)}.$$

Therefore

$$\rho_{F,\Psi} = \rho_{\{(f \circ g),\Psi\}} \le \max\left\{\rho_{(g_{e_1},\Psi)}, \rho_{(g_{e_2},\Psi)}\right\} = \rho_{g,\Psi},$$
  
i.e., $\rho_{F,\Psi} \le \rho_{g,\Psi}.$  (5)

Now using Lemma 3 and Theorem 1, we have

$$\rho_{(F_{e_1},\Psi)} = \rho_{\{(f_{e_1} \circ g_{e_1}),\Psi\}} = \limsup_{r_1 \to \infty} \frac{\log \log M_1(r_1^{1+\delta}, f_{e_1} \circ g_{e_1})}{\log \mathbb{E} \Psi(r_1^{1+\delta})}$$

$$\geq \liminf_{r_1 \to \infty} \frac{\log \log M(M(r_1, g_{e_1}), f_{e_1})}{\log \log M(r_1, g_{e_1})} \cdot \limsup_{r_1 \to \infty} \frac{\log \log M(r_1, g_{e_1})}{\log \mathbb{E} \Psi(r_1)]}$$

$$\geq \lambda^*_{(f_{e_1},\Psi)} \cdot \rho_{(g_{e_1},\Psi)} = 1 \cdot \rho_{(g_{e_1},\Psi)} = \rho_{(g_{e_1},\Psi)}$$

Similarly,  $\rho_{(F_{e_2}, \Psi)} \ge \rho_{(g_{e_2}, \Psi)}$ .

Therefore

$$\rho_{F,\Psi} = \rho_{\{(f \circ g),\Psi\}} \ge \max\left\{\rho_{(g_{e_1},\Psi)}, \rho_{(g_{e_2},\Psi)}\right\} = \rho_{g,\Psi},$$
  
i.e.,  $\rho_{F,\Psi} \ge \rho_{g,\Psi}.$  (6)

Therefore from (5) and (6), the result follows.

**Theorem 3.** Let f(w) and g(w) be any two bicomplex potential fluid flows satisfying the following properties :

(i) f(w) is entire and g(w) is transcendental such that  $\rho_{fog} = 0$  and  $\lambda_g < \infty$ .

Then  $\rho_{f,\Psi}^{**}\lambda_{g,\Psi}^{**} \le \rho_{(f \circ g),\Psi}^{**} \le \rho_{f,\Psi}^{**}\rho_{g,\Psi}^{**}$ .

**Proof.** Using Lemma 3 we have

$$\begin{split} \rho_{(F_{e_{1},\Psi})}^{**} &= \rho_{\{(f_{e_{1}} \circ g_{e_{1}}),\Psi\}}^{**} = \limsup_{r_{1} \to \infty} \frac{\log M_{1}(r_{1}^{1+\delta}, f_{e_{1}} \circ g_{e_{1}})}{\log [\Psi(r_{1}^{1+\delta})} \\ &\geq \limsup_{r_{1} \to \infty} \frac{\log M_{1}(M_{1}(r_{1}, g_{e_{1}}), f_{e_{1}})}{\log M_{1}(r_{1}, g)} \cdot \liminf_{r_{1} \to \infty} \frac{\log M_{1}(r_{1}, g_{e_{1}})}{\log [\Psi(r_{1})} \\ &= \rho_{f_{e_{1}}}^{**} \lambda_{g_{e_{1}}}^{**}. \end{split}$$

Similarly,

$$\rho_{(F_{e_2},\psi)}^{**} = \rho_{\{(f_{e_2} \circ g_{e_2}),\psi\}}^{**} \ge \rho_{(f_{e_2},\psi)}^{**} \cdot \lambda_{(g_{e_2},\psi)}^{**}$$

Therefore

$$\rho_{fog,\Psi}^{**} = max\left\{\rho_{\{(f_{e_1},\Psi)o(g_{e_1},\Psi)\}}^{**}, \rho_{\{(f_{e_2},\Psi)o(g_{e_2},\Psi)\}}^{**}\right\} \ge \rho_{f,\Psi}^{**}.\lambda_{g,\Psi}^{**}$$
(7)

Again, by using Lemma 2 we have

$$\begin{split} \rho_{(F_{e_1},\Psi)}^{**} &= \rho_{f_{e_1}og_{e_1},\Psi}^{**} = \limsup_{r_1 \to \infty} \frac{\log M_1(r_1^{1+\delta}, f_{e_1}og_{e_1})}{\log \Psi(r_1^{1+\delta})} \\ &\leq \limsup_{r_1 \to \infty} \frac{\log M_1(M_1(r_1, g_{e_1}), f_{e_1})}{\log M_1(r_1, g)} \cdot \limsup_{r_1 \to \infty} \frac{\log M_1(r_1, g_{e_1})}{\log [\Psi(r_1)]} \\ &= \rho_{f_{e_1}}^{**} \rho_{g_{e_1}}^{**}. \end{split}$$

Similarly,

$$\rho_{(F_{e_2},\psi)}^{**} = \rho_{\{(f_{e_2} \circ g_{e_2}),\psi\}}^{**} \le \rho_{(f_{e_2},\psi)}^{**} \cdot \rho_{(g_{e_2},\psi)}^{**}.$$

Therefore

$$\rho_{(fog),\psi}^{**} = \max\left\{\rho_{(f_{e_1},\psi)o(g_{e_1},\psi)}^{**}, \rho_{(f_{e_2},\psi)o(g_{e_2},\psi)}^{**}\right\} \le \rho_{f,\psi}^{**}\rho_{g,\psi}^{**}.$$
 (8)

Therefore, from (7) and (8) the result follows.

### **CONCLUSION AND FUTURE PROSPECTS:**

The estimates and derivation carried out in the paper may also be established under the treatment of bicomplex valued functions of slower and faster growth.

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